An asymptotic regime

Qualitative properties of solutions of the nonlinear Schrödinger equation on metric graphs Nonlinear Elliptic PDE in Hauts-de-France, 4th edition, Calais Damien Galant

CERAMATHS/DMATHS

Université Polytechnique Hauts-de-France Département de Mathématique

Université de Mons F.R.S.-FNRS Research Fellow



Joint work with Colette De Coster (CERAMATHS/DMATHS, Valenciennes, France) and Christophe Troestler (UMONS, Mons, Belgium)

Thanks to Colette for the slides!

Monday 16 June 2025

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Compact metric graphs

A compact metric graph $\mathcal{G} = (\mathbb{V}, \mathbb{E})$ is a connected network made up of a finite number of finite length edges $e \in \mathbb{E}$, glued at vertices $v \in \mathbb{V}$, according to the topology of a graph.



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Any bounded edge e is identified with a compact interval of ℝ;
u ∈ L^p(G) ⇔ u ∈ L^p(e) for every edge e of G.

The Sobolev space $H^1(\mathcal{G})$

The Sobolev space $H^1(\mathcal{G})$ is defined as follows

$$u \in H^1(\mathcal{G}) \iff egin{cases} u \in H^1(e) & ext{ for every edge } e ext{ of } \mathcal{G}, \ u : \mathcal{G} o \mathbb{R} & ext{ is continuous on } \mathcal{G}. \end{cases}$$

Here is what a typical $H^1(\mathcal{G})$ function looks like:



Introduction

An asymptotic regime

The differential system

Given constants p > 2 and $\lambda > 0$, we are interested in solutions $u \in L^2(\mathcal{G})$ of the differential system

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$$\begin{cases} -\tilde{u}'' + \lambda \tilde{u} = |\tilde{u}|^{p-2}\tilde{u} & \text{ on every edge of } \mathcal{G}, \\ \tilde{u} \text{ is continuous} & \text{ on } \mathcal{G}, \\ \sum_{e \succ V} \tilde{u}'_e(V) = 0 & \text{ for every } V \in \mathbb{V} \setminus Z, \\ \tilde{u}(V) = 0 & \text{ for every } V \in Z, \end{cases}$$

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$\begin{bmatrix} e \succ v \\ \tilde{u}(v) = 0 \end{bmatrix}$	for every $v \in Z$,

Here, Z is a set of degree-one vertices where we impose the homogenous Dirichlet boundary condition. For $v \in V \setminus Z$, the condition on the sum of derivatives is called *Kirchhoff's condition*.

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The nonlinear Schrödinger equation Kirchhoff's condition: degree-one nodes



An asymptotic regime

Examples

The nonlinear Schrödinger equation Kirchhoff's condition in general: outgoing derivatives



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Variational formulations

Solutions of our problem correspond to critical points of the action functional J_λ defined by

$$J_{\lambda}(u) := \frac{1}{2} \int_{\mathcal{G}} |u'|^2 \, \mathrm{d}x + \frac{\lambda}{2} \int_{\mathcal{G}} |u|^2 \, \mathrm{d}x - \frac{1}{p} \int_{\mathcal{G}} |u|^p \, \mathrm{d}x$$

on the Sobolev space

$$H^1_Z := \Big\{ u : \mathcal{G} \to \mathbb{R} \ \Big| \ u \text{ is continuous; } u, u' \in L^2(\mathcal{G}); \ \forall v \in Z, u(v) = 0 \Big\}.$$

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We are interested in the qualitative properties of

1) solutions minimizing the action on the set of nonzero solutions \rightarrow

the ground states (GS)

2) solutions minimizing the action on the set of nodal solutions \rightarrow

the nodal ground states (NGS)

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How do symmetries of the graph ${\mathcal G}$ transfer to symmetries of the GS and the NGS?

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How do symmetries of the graph ${\mathcal G}$ transfer to symmetries of the GS and the NGS?

Where are the roots of the NGS? The maximum value? ...

Introduction

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Example 1 – The segment with two points glued together



What is the shape of the GS? The NGS?

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The asymptotic regime $p \rightarrow 2$

Hope: obtain more information in the regime $p \approx 2$, by studying the *spectral* properties of the problem.

For every positive integer k and p > 2, we want to relate the solutions of the nonlinear problem

$\int -\tilde{u}'' + \lambda \tilde{u} = \tilde{u} ^{p-2}\tilde{u}$	on every edge of $\mathcal{G},$
\tilde{u} is continuous	on $\mathcal{G},$
$\sum_{e \in V} \tilde{u}'_e(v) = 0$	for every $v \in V \setminus Z$,
$\widetilde{u}(V) = 0$	for every $v \in Z$,

to the eigenfunctions of the corresponding eigenvalue problem with eigenvalue γ_k .

Introduction

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A rescaling

In order to better understand the behaviour of the solutions as $p \to 2$, we consider the new variable $u = \gamma_k^{-1/(p-2)} \tilde{u}$. They are solutions of the nonlinear problem

$$\begin{cases} -u'' + \lambda u = \gamma_k |u|^{p-2} u & \text{ on every edge of } \mathcal{G}, \\ u \text{ is continuous} & \text{ on } \mathcal{G}, \\ \sum_{e \succ v} u'_e(v) = 0 & \text{ for every } v \in \mathbb{V} \setminus Z, \\ u(v) = 0 & \text{ for every } v \in Z. \end{cases}$$
 $(\mathcal{P}_{p,k})$

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Let $(u_{p_n})_n$ be a sequence of solutions to $(\mathcal{P}_{p_n,k})$, $(p_n)_n \subseteq]2, +\infty[, p_n \rightarrow 2.$

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Assume that

$$u_{p_n} \xrightarrow[n \to \infty]{H^1_Z} u_*$$

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Question

What can we say about u_* ?

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Therefore, u_* belongs to E_k .

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Question

Is that all we can say about u_* ?

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Let us use specifically $\psi \in E_k$ as a test function in $(\mathcal{P}_{p_n,k})$. We obtain

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Using u_{p_n} as a test function in the equation $-\psi'' + a\psi = \lambda_k \psi$, we get

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Thus,

$$\int_{\mathcal{G}} (|u_{p_n}|^{p_n-2}-1) u_{p_n} \psi \, \mathrm{d} x = 0.$$

We divide by $p_n - 2$:

$$\int_{\mathcal{G}} \frac{|u_{p_n}|^{p_n-2}-1}{p_n-2} u_{p_n} \psi \, \mathrm{d} x = \int_{\mathcal{G}} \frac{\mathrm{e}^{(p_n-2)\ln|u_{p_n}|}-1}{p_n-2} u_{p_n} \psi \, \mathrm{d} x = 0.$$

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Taking $n \to \infty$ leads to

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Taking $n \to \infty$ leads to

$$\int_{\mathcal{G}} (u_* \ln |u_*|) \psi \, \mathrm{d}x = 0.$$

Definition

A function $u_* \in E_k$ is a **solution of the reduced problem on** E_k if and only if

$$\int_{\mathcal{G}} (u_* \ln |u_*|) \psi \, \mathrm{d}x = 0$$

for all $\psi \in E_k$.

Introduction	An asymptotic regime	Examples

Recap

Given a sequence $(u_{p_n})_n$, $p_n \to 2$ converging weakly to $u_* \in H^1_Z$, we have seen that necessarily:

Introduction	

Recap

Given a sequence $(u_{p_n})_n$, $p_n \to 2$ converging weakly to $u_* \in H^1_Z$, we have seen that necessarily:

• u_* belongs to E_k ;

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Recap

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- u_* belongs to E_k ;
- u_* is a solution of the reduced problem, namely be so that

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Recap

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for all $\psi \in E_k$.

Question

Given a solution of the reduced problem $u_* \in E_k$, can one find solutions of $(\mathcal{P}_{p,k})$ close to u_* for $p \approx 2$? Can one detect when there is only one solution close to u_* for a given $p \approx 2$?

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Functional space with extra regularity:

$$H := \left\{ u \in H^1_Z \mid u \text{ is } H^2 \text{ in each edge}, u \text{ satisfies Kirchhoff's conditions} \right\}.$$

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Functional space with extra regularity:

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We fix $k \geq 1$ and we define the map

$$F:\begin{cases} [2,+\infty[\times H \to L^2(\mathcal{G}),\\ (p,u) \mapsto -u'' + \lambda u - \lambda_k |u|^{p-2}u. \end{cases}$$

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and when p > 2,

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Examples

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we obtain good invertibility properties on E_k^{\perp} and we are then reduced to a finite dimensional problem on E_k .

Examples

A word of caution



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Examples

A word of caution

Be careful! A Implicit Function Theorems require regularity!

To perform the Lyapunov-Schmidt reduction around u_* , we will need

$$F:\begin{cases} [2,+\infty[\times H \to L^2(\mathcal{G}),\\ (p,u) \mapsto -u'' + \lambda u - \lambda_k |u|^{p-2}u. \end{cases}$$

to be C^2 in u in the neighborhood of $(2, u_*)$.

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Examples

An important set

Expressions such as $u\mapsto u\,\ln|u|$ and its derivative $u\mapsto 1+\ln|u|$ appear in the study.

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Definition (An important set)

$$S := \Big\{ u \in H \mid \inf_{x \in \mathcal{G}} (|u(x)| + |u'(x)|) > 0 \Big\}.$$

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Remark: if $u \in E_k$, then

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On graphs, this is not automatic: no unique continuation principle!

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Nondegenerate solutions of the reduced problem

Definition

A solution $u_* \in E_k \cap S$ of the reduced problem on E_k is **nondegenerate** if and only if the map

$$E_k o E_k : \psi \mapsto P_{E_k} ((1 + \ln |u_*|)\psi)$$

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Remark: nondegeneracy always holds if dim $E_k = 1$.

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Examples

Main Theorem

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- I non-existence: If u_∗ is not a solution of the reduced problem, then there exists a neighbourhood U of (2, u_∗) in [2, +∞[× H so that problem (P_{p,k}) has no solution in U with p > 2;
- 2 existence, uniqueness and non-degeneracy: If u_{*} is a nondegenerate solution of the reduced problem, then there exists a neighbourhood U of (2, u_{*}) in [2, +∞[× H and a number ε > 0 so that for all p ∈]2, 2 + ε], there exists a unique u_p ∈ H so that (p, u_p) belongs to U and so that u_p is a solution of problem (P_{p,k}).

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Unidimensional eigenspaces

In case $E_k = \operatorname{span} \varphi$ is of dimension 1, up to sign, we know exactly the limit u^* as we know that $u_* = a\varphi$ with a such that

$$0 = \int_{\mathcal{G}} \varphi^2 \ln |a\varphi| \, \mathrm{d}x = \int_{\mathcal{G}} \varphi^2 (\ln |a| + \ln |\varphi|) \, \mathrm{d}x.$$

Moreover, it is nondegenerate.

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Theorem

If $p \approx 2$ is close enough to 2, the positive solution of $(\mathcal{P}_{p,1})$ is unique and is a ground state of the problem.

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Show that there exists C > 0 such that all positive solutions of (P_{p,1}) with 2 H¹(G)</sub> ≤ C;

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- When p → 2, sequences of positive solutions to (P_{p,1}) converge weakly (up to subsequences) to the only positive eigenfunction;
- Since dim $E_1 = 1$, u_* is nondegenerate;
- The Lyapunov-Schmidt reduction proves the uniqueness result.

Convergence of nodal ground states when $p \rightarrow 2$

Theorem (Convergence of nodal ground states)

If $(u_{p_n})_n$ is a sequence of nodal ground states of $(\mathcal{P}_{p,k})$ with $p_n \to 2$, then up to a subsequence one has that

$$u_{p_n} \xrightarrow[n \to \infty]{H^2} u_*,$$

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Convergence of nodal ground states when p ightarrow 2

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where $u_* \in E_2 \setminus \{0\}$ is a solution of the reduced problem.

Remark

If u_* belongs to S (i.e. does not vanish on any edge) and is nondegenerate, one may obtain uniqueness and symmetry results by the Lyapunov-Schmidt reduction.

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The *n*-bridge



n-edges e_1 , ... e_n of length $2L_1$, ..., $2L_n$ with $L_1 > L_2 \ge L_3 \ge ... \ge L_n$. What can be said on the ground state and the nodal ground state for $p \approx 2$?

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Ground state: easy ... It is constant for $p \approx 2$. What about the nodal ground state?

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The *n*-bridge The second eigenspace

Let us parametrize the edges on $[-L_i, L_i]$. The solution of

 $\begin{cases} -u'' = \gamma u & \text{on } [-L_i, L_i], \\ u \text{ is continuous } & \text{on } \mathcal{G}, \\ \sum_{e \succ V} \tilde{u}'_e(V) = 0 & \text{ for every } V \in \mathbb{V} \end{cases}$

are given by $u_i(x) = a_i \cos(\sqrt{\gamma}x) + b_i \sin(\sqrt{\gamma}x)$ with, for all $1 \le i,j \le n$,

$$\begin{aligned} a_i \cos(\sqrt{\gamma}L_i) &= a_j \cos(\sqrt{\gamma}L_j), \\ b_i \sin(\sqrt{\gamma}L_i) &= b_j \sin(\sqrt{\gamma}L_j), \\ \sum a_i \sin(\sqrt{\gamma}L_i) &= 0, \\ \sum b_i \cos(\sqrt{\gamma}L_i) &= 0. \end{aligned}$$

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The *n*-bridge The second eigenspace

We prove that the second eigenvalue is defined by

$$\gamma_2 \in \left] \left(\frac{\pi}{2L_1} \right)^2, \ \min\left\{ \left(\frac{\pi}{2L_2} \right)^2, \left(\frac{\pi}{L_1} \right)^2 \right\} \right[\ \text{is the solution of} \ \sum \tan(\sqrt{\gamma_2}L_i) = 0,$$

with eigenfunction

$$\begin{cases} \varphi_{2,1}(x) &= a_1 \cos(\sqrt{\gamma_2} x), \\ \varphi_{2,i}(x) &= a_1 \frac{\cos(\sqrt{\gamma_2} L_1)}{\cos(\sqrt{\gamma_2} L_i)} \cos(\sqrt{\gamma_2} x), \quad \text{for } i \ge 2 \end{cases}$$

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The *n*-bridge Properties of φ_2

We observe that:

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The *n*-bridge Properties of φ_2

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The *n*-bridge Properties of φ_2

We observe that:

- **1** $\varphi_{2,i}$ are even on each edge;
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$$A_i = \left|a_1 \frac{\cos(\sqrt{\gamma_2}L_1)}{\cos(\sqrt{\gamma_2}L_i)}\right|$$
 then $A_1 > A_2 \ge \ldots \ge A_n$;

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Ground state: It is constant for $p \approx 2$.

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Ground state: It is constant for $p \approx 2$.

Nodal ground state: For $p \approx 2$

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The *n*-bridge Conclusions

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Nodal ground state: For $p \approx 2$

1 u_i are even on each edge;

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The *n*-bridge Conclusions

Ground state: It is constant for $p \approx 2$.

Nodal ground state: For $p \approx 2$

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The *n*-bridge Conclusions

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The *n*-bridge Conclusions

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5 if
$$L_i = L_j$$
 then $u_i(x) = u_j(x)$.

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The dumbbell – Conclusion



Ground state: It is constant for $p \approx 2$.

Nodal ground state: For $p \approx 2$

- 1 u_1 , u_2 are even,
- **2** the root of u is the middle point of e_3 ,
- <u>3</u> *u* is odd "globally".
- 4 u_3 is strictly monotone.
- 5 the maximum of φ_2 is the "extremity" of the loop.

Examples

The tadpole – Dirichlet



For $p \approx 2$, the positive GS is even on the loop, increasing on the segment.

For $p \approx 2$, the NGS is

- even on the loop,
- 2 one nodal domain is included in the loop,
- 3 the maximum of the amplitude is in the interior of the segment.

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The tadpole – Neumann



For $p \approx 2$, the GS is constant

The second eigenvalue is $\gamma_2 = (\frac{\pi}{2L})^2$ with eigenfunction

$$\varphi_{2,1}(x) = -2a_2\cos\left(\frac{\pi}{2L}x\right), \quad \varphi_{2,2}(x) = a_2\cos\left(\frac{\pi}{2L}x\right)$$

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The tadpole – Neumann



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$$\varphi_{2,1}(x) = -2a_2\cos\left(\frac{\pi}{2L}x\right), \quad \varphi_{2,2}(x) = a_2\cos\left(\frac{\pi}{2L}x\right)$$

1 φ_2 is even on the loop,

- 2 the loop is one nodal domain, the segment is the second nodal domain, the nodal set is the vertex v
- 3 the maximum of the amplitude is on the vertex of degree 1.
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Examples

The tadpole – Neumann



What about the NGS for $p \approx 2$?

Easy:

- 1 *u* is even on the loop,
- 2 the maximum of the amplitude is on the line.

Examples

The tadpole - Neumann



What about the nodal domain? u(v) = 0 or not?

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Examples

The tadpole - Neumann

u cannot be equal to zero at the vertex v as otherwise u_2 is a solution of

$$\begin{cases} -u'' + \lambda u = |u|^{p-2}u \\ u(-L) = u(L) = 0 \\ u > 0 \text{ on }] - L, L[\end{cases}$$

and u_1 is a solution of

$$\begin{cases} -u'' + \lambda u = |u|^{p-2}u \\ u'(0) = u(L) = 0 \\ u < 0 \text{ on }]0, L[\end{cases}$$

By uniqueness of the solution of these problems and as $u_1 = -u_2|_{[0,L]}$ this is not a solution on the graph as it does not satisfy the Kirchhoff condition.

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The tadpole – Neumann



In fact $u^{-1}(0) = \{x_0\}$ with x_0 a point of the segment.

$$\varphi_{2,1}(x) = -2a_2\cos(\frac{\pi}{2L}x), \quad \varphi_{2,2}(x) = a_2\cos(\frac{\pi}{2L}x)$$

hence the amplitude is larger on the segment.

The same is true for the NGS by convergence.

The tadpole – Neumann

We know that the time needed to go from the maximum to 0 is a decreasing function of the value of the maximum. Hence the result.

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The tadpole – Neumann

We know that the time needed to go from the maximum to 0 is a decreasing function of the value of the maximum. Hence the result.

Conclusion: For $p \approx 2$, the NGS satisfies :

- 1 *u* is even on the loop,
- 2 the maximum of the amplitude is on the line,
- 3 $u^{-1}(0) = \{x_0\}$ with x_0 a point of the segment,
- 4 one nodal domain is included in the segment, the other contains the loop.

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One bubble



One bubble

Second eigenvalue:

If $L_3 < 4L_1$ then $\gamma_2 = \left(\frac{3\pi}{2(L_3 + 2L_1)}\right)^2$ is simple with the second

eigenfunction even on the loop and not identically zero on an edge. In that case, the NGS is also even on the loop.

One bubble

Second eigenvalue:

If $L_3 < 4L_1$ then $\gamma_2 = \left(\frac{3\pi}{2(L_3 + 2L_1)}\right)^2$ is simple with the second eigenfunction even on the loop and not identically zero on an edge. In that case, the NGS is also even on the loop. If $L_3 > 4L_1$ then $\gamma_2 = \left(\frac{\pi}{L_3}\right)^2$ is simple with the second eigenfunction odd on the loop and identically zero on e_1 and on e_2 . What about the NGS?

Examples

The limit when $p \rightarrow 2$ on the loop



Examples

Box when $p \approx 2$ on the loop



Examples

Sign change on the loop



Problem: behaviour at the node?



Introduction

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Continuity



Examples

Conclusion: On the loop



Introduction

One bubble – Asymmetric vertex conditions



The GS is even on the loop.

If $L_3 < 4L_1$ then the NGS is also even on the loop.

If $L_3 > 4L_1$ then the NGS is odd on the loop and identically zero on e_1 and on e_2 .

Introduction	

The symmetric stargraph – Symmetry breaking



For *L* fixed, by uniqueness, for $p \approx 2$, the GS is *symmetric* (i.e. its restrictions to all edges, viewed as functions $[0, L] \rightarrow \mathbb{R}$, are all equal).

The symmetric stargraph – Symmetry breaking



For *L* fixed, by uniqueness, for $p \approx 2$, the GS is *symmetric* (i.e. its restrictions to all edges, viewed as functions $[0, L] \rightarrow \mathbb{R}$, are all equal). Instead, for any p > 2, if *L* is large enough then the ground state on \mathcal{G}_L is

not symmetric.

The symmetric stargraph – Symmetry breaking



For *L* fixed, by uniqueness, for $p \approx 2$, the GS is *symmetric* (i.e. its restrictions to all edges, viewed as functions $[0, L] \rightarrow \mathbb{R}$, are all equal).

Instead, for any p > 2, if *L* is large enough then the ground state on G_L is *not* symmetric.

In particular, the uniqueness of the positive solution is not always valid (not as on the interval).

Thanks for your attention!

Qualitative properties for NLS on compact graphs

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